

Equation (8) is therefore a closed-form coupled twist distribution for a composite wing with the warping effects accounted for, where

$$\bar{y}_0 = y_0/\ell_0 \quad (9)$$

When Eq. (8) is evaluated at the wing tip and compared to an equivalent expression predicted by St. Venant's theory, the following expression is obtained:

$$\frac{\alpha_0(1)}{\alpha_0(1)_{\text{St.V}}} = 1 - \frac{\tanh 4\bar{\lambda}_c}{2\bar{\lambda}_c} - \frac{1}{9\bar{\lambda}_c^2} \left(\frac{1}{\cosh 4\bar{\lambda}_c} - 1 \right) \quad (10)$$

where $\alpha_0(1)$ is the wing tip twist given by Eq. (8) while $\alpha_0(1)$ is the wing tip twist given by the St. Venant torsion theory. A plot of Eq. (10) is shown in Fig. 1.

Steady-State Concentrated Tip Twist Loads

If the wing is under the influence of a concentrated twisting moment F_0 at the tip as a result of a steady-state coupled bending torsion displacement, the exact closed-form twist distribution that satisfies these equations of motion and their associated boundary conditions is given by

$$\alpha_0(\bar{y}_0) = \frac{6F_0 \ell_0}{c_0^3(4\bar{\lambda}_c)^2} \left[\bar{y}_0 - \frac{\sinh 4\bar{\lambda}_c \bar{y}_0}{4\bar{\lambda}_c} + \frac{\tanh 4\bar{\lambda}_c}{4\bar{\lambda}_c} \times (\cosh \bar{\lambda}_c \bar{y}_0 - 1) \right] \quad (11)$$

When the twist distribution given by Eq. (11) is evaluated at the wing tip and compared to its counterpart predicted by the St. Venant torsion theory, the following expression is obtained:

$$\frac{\alpha_0(1)}{\alpha_0(1)_{\text{St.V}}} = 1 - \frac{\tanh 4\bar{\lambda}_c}{4\bar{\lambda}_c} \quad (12)$$

It should be noted that the ratio given by Eq. (12) was plotted for the real values of $\bar{\lambda}_c$ in Refs. 5 and 6 and was shown to represent conditions where any errors resulting from using St. Venant's torsion theory are conservative (overdesign rather than underdesign). In this analysis, Eq. (12) is examined when $\bar{\lambda}_c$ is imaginary, which is possible if L_2 is very large. Under such circumstances, Eq. (12) becomes

$$\frac{\alpha_0(1)}{\alpha_0(1)_{\text{St.V}}} = 1 - \frac{\tan 4\bar{\lambda}_c}{4\bar{\lambda}_c} \quad (13)$$

Figure 2 depicts the conditions given by Eq. (13). It is therefore seen from the figure that there are certain ranges of $\bar{\lambda}_c$ for which nonconservative errors are possible by using the St. Venant twist theory.

Results and Conclusions

The results are shown in Figs. 1 and 2. Figure 1 shows a comparison of the static wing tip twist obtained in the present study and that obtained via St. Venant's twist theory in the presence of statically distributed forces and low-to-moderate coupling. Figure 2 shows the trend for concentrated forces and substantial coupling. In Fig. 1, it is seen that the presence of coupling makes the errors of St. Venant's theory worse. This seems to suggest that the more sophisticated theory is more important for wings with coupling (e.g., wings aeroelastically tailored using elastic cross coupling).

Figure 2 also shows that nonconservative errors $\{|\alpha_0(1)/[\alpha_0(1)_{\text{St.V}}]| > 1\}$ are possible.

Using Figs. 1 and 2, the following conclusions can be summarized: 1) ignoring warping arbitrarily using St. Venant's theory could result in very significant errors (as high as over 80%) in analytical results for composite aircraft wings; 2) warping is more important (St. Venant's theory is less accurate)

rate) for wings with coupling; and 3) St. Venant's theory (which has always been shown to be conservative^{1,2} can be nonconservative or St. Venant's approximation can lead to an unsafe design error (underdesign rather than overdesign from a stability point of view).

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Investigation of Internal Singularity Methods for Multielement Airfoils

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I. Introduction

A NUMBER of integral equation methods are currently used for solving linear potential flows about single airfoils, e.g., Hess and Smith,¹ Bristow,² Hess,^{3,4} Maskew and Woodward,⁵ and Chen and Dalton,⁶ and about multielement airfoils, e.g., Banyopadhyay et al.⁷ The comparisons of the solutions of these numerical methods for single airfoils are shown in Ref. 8. The analytical solution for two-dimensional potential flow past the airfoils had been presented by Williams⁹ and James.¹⁰ The method developed by Garrick¹¹ can analyze flow about two-element airfoils. The method developed by Ives,¹² Halsey,¹³ and Suddhoo and Hall¹⁴ can analyze the flow about multielement airfoils.

A numerical solution, which is obtained by the internal linear-vortex-source flat-element method, is compared with the

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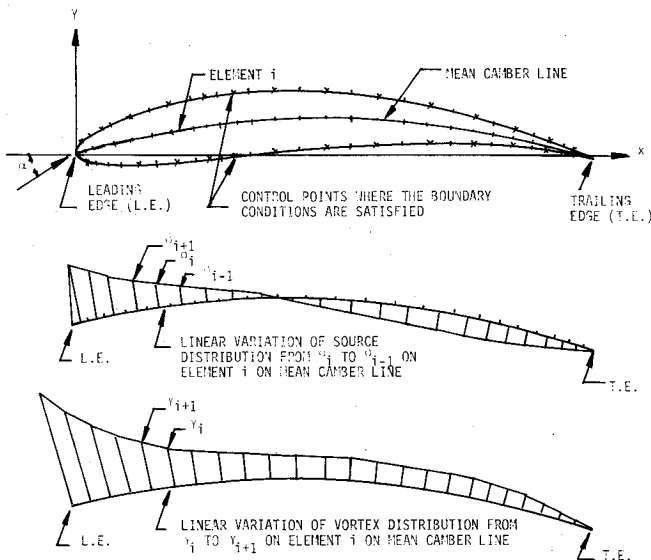


Fig. 1 Distribution of singularities and control points for internal linear vortex-source model.

analytical solution, which can be determined by conformal transformation.¹⁵ The internal integral equation method developed by Chen and Dalton⁶ had the same numerical instability of the pressure distributions on the airfoil surface as that developed by Banyopadhyay et al.⁷ The solutions of this numerical method are compared with those of the internal singularity method that was developed by Banyopadhyay et al.⁷

The modification of this method by using the curved element modeling of Hess³ is applied to the multielement airfoils in order to investigate the effects of the curved elements on the distributions of the pressure and the strength of the singularity of this numerical method.

II. Analysis of Internal Singularity Method

This numerical method was developed and is shown in Fig. 1. The mean camber line is divided into N elements from the leading edge to the trailing edge (N unknown linear variations of source strengths are placed on the midpoint of each element, and $N + 1$ unknown linear variations of vorticity strengths are placed on both ends of each element). Distributions of the vorticity and source strengths vary linearly across each element on the mean camberline inside the airfoil. These $(2N + 1)$ unknowns are determined by satisfying the boundary condition of tangency of flow at $2N$ control points on the airfoil surface together with appropriate Kutta condition at the trailing edge. Corresponding to the midpoint of each element on the mean camber line, two points can be defined on the upper and the lower airfoil surfaces from the intersection of normals drawn from the mean camber line with the airfoil profile (Fig. 1). This identifies $2N$ control points on the airfoil surface where the tangency boundary conditions are to be satisfied. The Kutta condition is satisfied by equating the tangential velocities at the control points adjacent to the trailing edge on the upper and the lower airfoil surfaces.

In the methods that were developed in Refs. 6 and 7, a small gap was maintained between the leading edge and the element on the mean camber line closest to the leading edge of the airfoil in order to avoid infinite velocity near the leading edge. The size of the gap needed to be chosen carefully (see Refs. 6 and 7) or it would cause instability of the pressure distributions on the airfoil surface. This method does not have such disadvantage, i.e., there is no gap to be left between the leading edge and the element on the mean camberline closest to the leading edge. The effect of the higher-order approximation is investigated by a modification of this method with the application of curved element modeling of Hess.³ (The airfoil

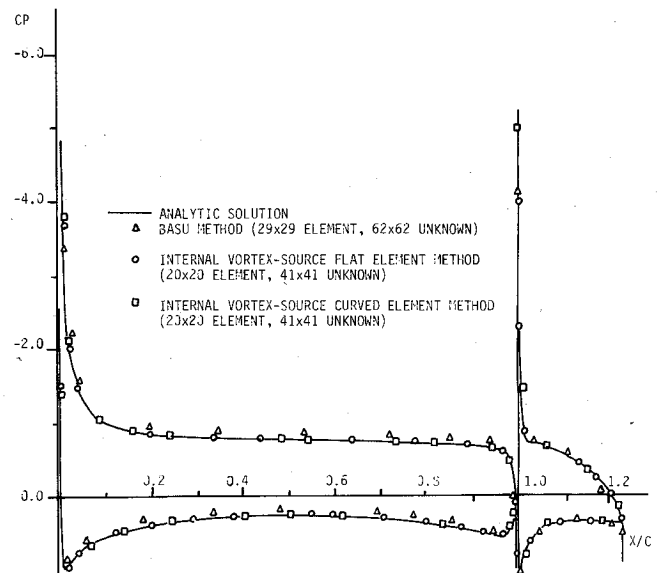


Fig. 2 Pressure distributions of internal singularity method for two-element airfoil.

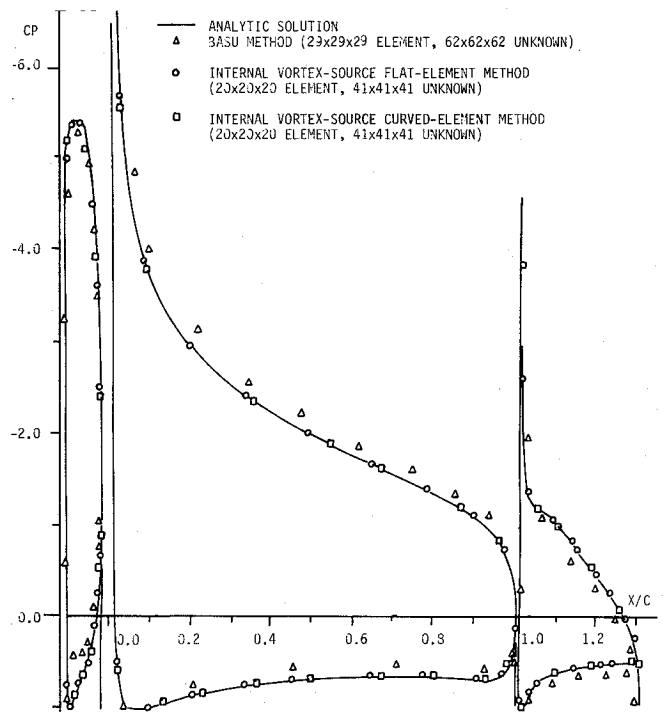


Fig. 3 Pressure distributions of internal singularity methods for three-element airfoil.

surface and the mean camberline can be represented by a number of curved elements, and the distributions of the vorticity and source strengths are parabolically varying across each curved element on the mean camber line.)

The conformal transformation of multiple circles onto multielement airfoils is a modification of the method developed by Suddhoo and Hall,¹⁴ i.e., the general Kármán-Trefftz transformation is used, and there are no coordinate translation and rotation in the calculation.

III. Results and Discussions

The pressure distributions for three internal singularity methods on the two-element airfoil are shown in Fig. 2. The agreement between the analytic solution and the numerical solutions for both the internal linear-vortex-source flat element and curved element methods is excellent. The pressure distributions on the first element calculated by the method of Ref. 7 is a little greater than the analytic solution.

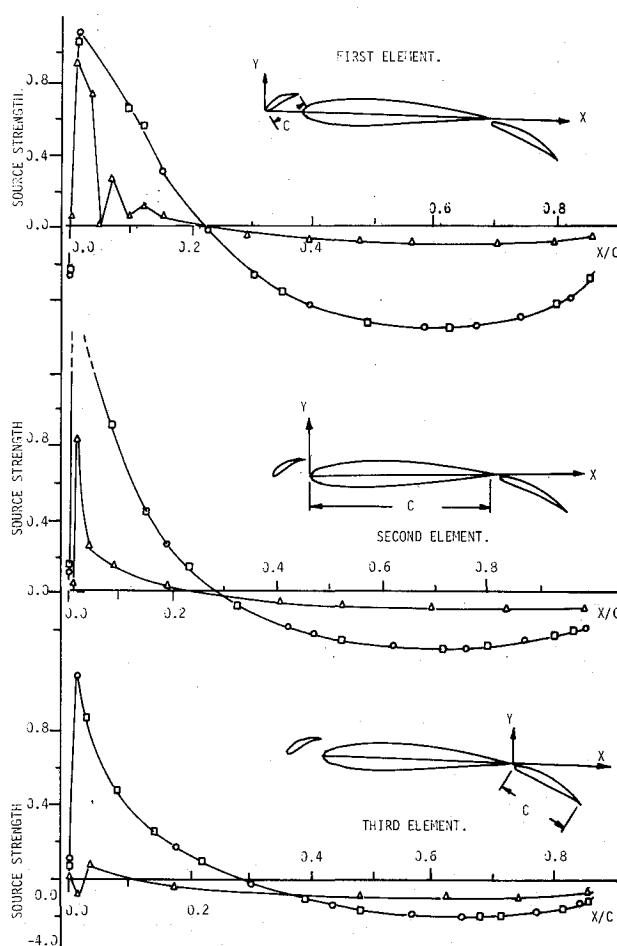
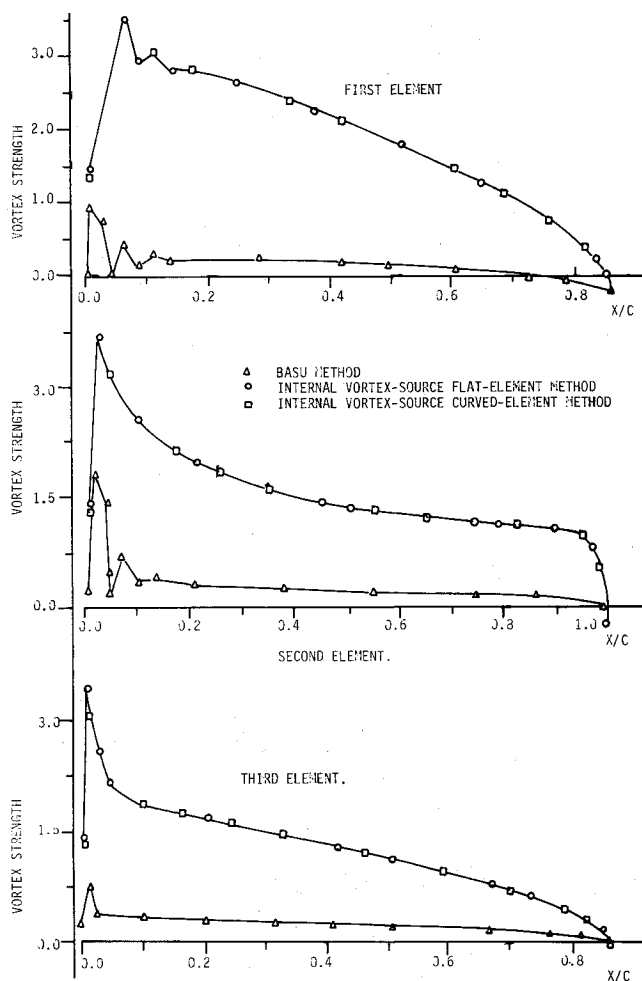


Fig. 4 Distributions of internal singularities of three-element airfoil.

The pressure distributions on the three-element airfoil are shown in Fig. 3. The results calculated by the method of Ref. 7 are very poor for the second and third airfoil elements when compared with the analytic solution. The agreement between the analytic solution and the solutions of both the internal linear-vortex-source flat element and curved element methods is good.

In order to investigate the discrepancy between the analytical solution and the numerical solution, which is calculated by the method developed by Ref. 7, the distributions of the source and vorticity strengths on the mean camber line of the three-element airfoil are shown in Fig. 4. The agreement of the distributions of the source and the vorticity strengths between the internal linear-vortex-source flat element and curved element methods is good. However, the distributions of the strengths of the singularities for the method of Ref. 7 are unstable in the region close to the nose of each airfoil element. The instability of the distributions of the singularities on the mean camber line of the three-element airfoil for the method of Ref. 7 causes the poor results of the pressure distributions that are shown in Fig. 3. It is suspected that due to the interferences of the source and the vorticity distributions on the mean camber line among the three airfoil elements, the criterion of the gap determination may need severe restriction for the method of Ref. 7.

IV. Conclusions

A flat element method based on linear distributions of vorticity and source on the mean camber line of the airfoil elements offers an approach to the solution of general

potential flow problems. This method does not need any gap to be left between the leading edge and the element on the mean camber line closest to the leading edge to avoid the instability of the numerical solutions. In view of the good agreement and the accuracy of the pressure distributions obtained by the flat element and the curved element methods, the higher-order approximation with curved element of this numerical method may be unnecessary in the calculation.

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